

文章编号:1009-9603(2022)06-0085-10

DOI:10.13673/j.cnki.cn37-1359/te.202107002

井洞相连的串珠状缝洞型油藏试井分析方法

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摘要:顺北油田包含溶洞、高角度裂缝、溶蚀孔洞和岩石基质4种储集介质,溶洞与高角度裂缝串联形成垂向串珠状缝洞体是常见的缝洞连通模式,钻井过程中经常出现钻柱放空和泥浆漏失的现象,这些复杂的地质和工程因素导致常规的多重介质模型难以拟合实测的试井数据。提出多个溶洞区域和裂缝区域串联的物理模型表征复杂的垂向串珠结构,溶洞区域仅包含溶洞介质,流体遵守自由流,裂缝区域包含高角度裂缝、溶蚀孔洞和基质,流体遵守达西定律,进而建立井洞相连的串珠状缝洞体试井模型,绘制试井典型图版,进行参数敏感性分析。结果表明:溶洞储集阶段和线性流阶段是串珠状缝洞体的典型流动阶段;根据试井压力导数曲线的斜率能识别溶洞和常规多重介质;重力影响下试井典型曲线会出现类似于封闭边界的压力响应特征,应采取措施降低重力的影响以提高串珠状缝洞体的开发效果。最后,应用所建模型解释了顺北油田典型井,除能解释常规的储层参数外,还能反演现有试井方法无法解释的溶洞等效半径、溶洞等效高度等参数。

关键词:溶洞;垂向串珠状缝洞体;重力;封闭边界;顺北油田

中图分类号:TE353

文献标识码:A

Well test analysis method for fracture-cavity reservoirs of beads-on-string structure with wellbore-cave connection

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Abstract: Shunbei Oilfield contains four types of storage media (i.e., caves, high-angle fractures, vugs, and rock matrix). The vertical beads-on-string structure formed by caves and high-angle fractures in series is a common fracture-cavity connection pattern. During drilling, drilling break and mud loss are prone to occur. These complex geological and engineering factors make it difficult for traditional multiple media models to fit the measured well test data. This paper proposes a physical model with multiple cave regions and multiple fracture regions in connection in series to characterize the complex vertical beads-on-string structure. The cave regions only contain caves in which the fluids are assumed to obey free flow. The fracture regions contain high-angle fractures, vugs, and rock matrix in which the flow obeys Darcy's law. On this basis, the paper constructs a well test model for the beads-on-string structure with a wellbore-cave connection, draws type curves, and conducts a parameter sensitivity analysis. The following results are obtained: the cave storage regime and linear flow regime are the typical flow stages of vertical beads-on-string structure; the caves and conventional multiple media can be distinguished by the slope of pressure derivative curves; the type curves under the influence of gravity show pressure response characteristics similar to a closed boundary, so the measures should be taken to reduce the influence of gravity for improving the development effect of the vertical beads-on-string structure. Finally, the proposed model is applied to interpret the representative wells in Shunbei Oilfield. Besides the conventional reservoir parameters, the parameters that cannot be interpreted by existing well test methods can also be obtained by inversion, such as equivalent cave radius and equivalent cave

收稿日期:2021-07-11。

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height.

Key words: cave; vertical beads-on-string structure; gravity; closed boundary; Shunbei Oilfield

缝洞型油藏一直是油气田勘探开发的热点,近年来,中国在塔里木盆地取得了较好的勘探成果,陆续开发了塔河油田、顺北油田等多个大型缝洞型油藏^[1-2]。顺北油田是由大型构造断裂活动和多期岩溶作用形成的断控岩溶油藏,储层埋深超过6 km,地面露头、测井、岩心和地震剖面显示储层发育溶洞、高角度裂缝、溶蚀孔洞和岩石基质4种介质^[3]。溶洞和高角度裂缝的连通关系复杂,水平串珠结构和垂向串珠结构是油藏中最常见的缝洞连通模式^[4-5]。钻井过程中,钻杆断裂情况占23%,断裂最长可达29 m;泥浆漏失情况占63.7%,漏失量最大可达3 483.55 m³。这些复杂的地质情况和工程因素给储层评价带来了困难。

试井分析是评价储层的重要手段^[6-7],在 Warren-Root 模型^[8]、Kazemi 模型^[9]和 De-swaan 模型^[10]的基础上,中外学者提出了大量缝洞型油藏试井模型。然而,传统的模型一般考虑三重介质,模型的建立基于经典的达西定律,很难精确描述溶洞中流体流动,特别是当井直接钻入溶洞内^[11-15]。随着缝洞型油藏的开发,学者们认识到了溶洞的存在。常学军等考虑裂缝和溶洞与井筒相连,提出了三重介质油藏试井分析方法^[16],张福祥等建立了井钻入溶洞内的缝洞型油藏试井模型^[17],林加恩等使用管流与渗流耦合来解决溶洞内的渗流问题^[18],尹洪军等

将溶洞考虑为扩大的井筒,建立了溶洞发育的缝洞型油藏试井解释模型^[19],杜鑫等联立质量、能量和动量守恒方程,建立了波动和流动耦合模型来分析井底的压力动态^[20-21],在文献[21]的基础上,徐燕东分析了重力对试井特征曲线的影响^[22]。但是,上述研究仅考虑了一个溶洞的情况,而串珠状溶洞往往是多个溶洞连接在一起。

围绕串珠状溶洞的试井模型和渗流理论研究,LI等建立了横向串珠缝洞体产量递减分析模型^[23],DU等利用多区复合模型建立了横向串珠缝洞体试井模型^[24]。然而,上述方法只考虑了溶洞、裂缝和基质3种介质,没有考虑溶蚀孔洞;且没有提出针对垂向串珠结构的试井分析方法。

综合分析油藏的地质情况和工程因素,考虑4种介质和重力影响,建立井洞相连的垂向串珠结构试井模型,分析典型特征曲线,将提出的模型在顺北油田进行实例应用和试井解释,提供一种评估串珠状缝洞体储层参数的试井分析方法。

1 模型建立

1.1 物理模型

提出2个溶洞区域和裂缝区域串联的物理模型(图1a)来表征垂向串珠结构(图1b)。井筒直接和

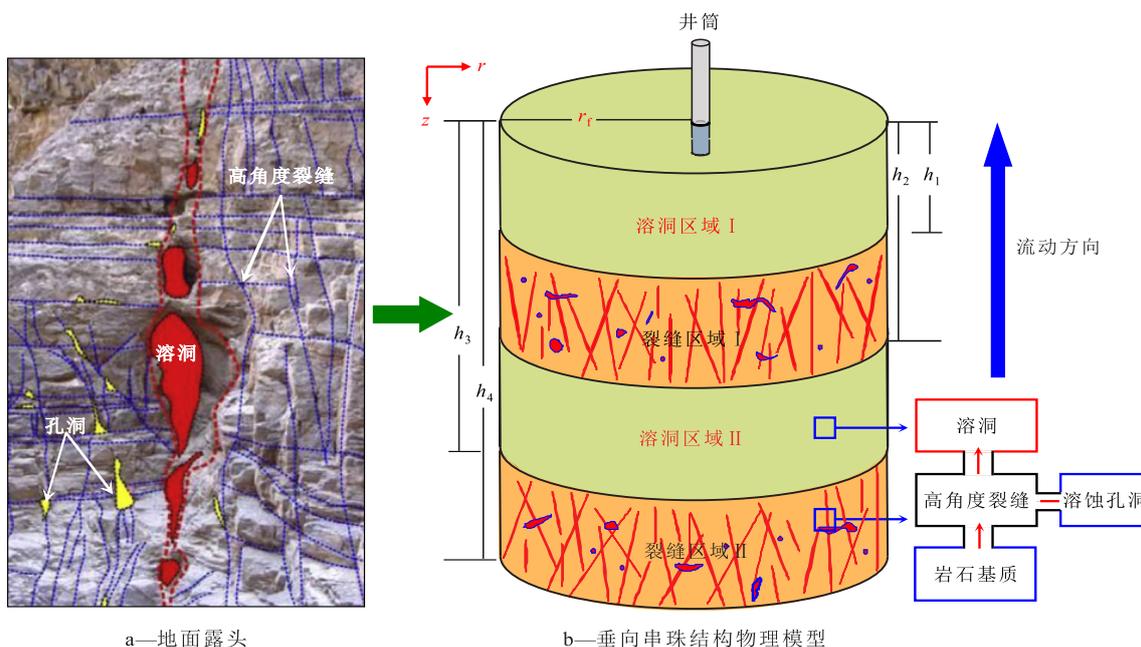


图1 垂向串珠结构示意图

Fig.1 Vertical beads-on-string structure

溶洞区域相连,溶洞区域仅包含溶洞介质,等效半径为 r_i , 溶洞区域 $i(i=1, 2)$ 的顶高和底高分别为 $h_{2(i-1)}$ 和 h_{2i-1} , 由于溶洞区域未充填, 流体假设为自由流(即流体穿过溶洞区域的压降很小, 建模时可以忽略); 裂缝区域包含高角度裂缝、溶蚀孔洞和基质介质, 裂缝区域 $j(j=1, 2)$ 的顶高和底高分别为 h_{2j-1} 和 h_{2j} , 由于裂缝区域内高角度裂缝比较发育, 流动方向考虑为沿垂直方向单向流动, 孔洞和基质内流体以拟稳态窜流方式流入裂缝。其他基本假设如下: ① 储层原始地层压力为 p_i , 单相微可压缩原油的体积系数为 B , 油井以定产量 q 生产。② 溶洞区域 i 的储集系数为 C_{vi} , 井筒储集系数为 C , 油藏的综合压缩系数为 c_i , 流体压缩系数为 c_{fi} 。③ 裂缝区域 j 中, 裂缝、孔洞和基质的渗透率分别为 K_{fj}, K_{vj}, K_{mj} , 孔隙度分别为 $\phi_{fj}, \phi_{vj}, \phi_{mj}$, 流体黏度分别为 $\mu_{fj}, \mu_{vj}, \mu_{mj}$ 。④ 考虑重力影响, 忽略毛管压力的影响。

1.2 数学模型

1.2.1 模型建立

根据基本假设, 溶洞区域内流体遵守自由流, 流体流过溶洞区域造成的压降可忽略, 易得溶洞区域 I 界面处的压力连续方程为:

$$p_{wf}(z=0, t) = p_{f1}(z=h_1, t) = p_{cv1} \quad (1)$$

考虑造成溶洞区域 I 内流体质量发生改变的 3 个因素: ① 溶洞区域 I 内流体流入井筒。② 裂缝区域 I 的流体流入溶洞区域 I。③ 岩石和流体的压缩性。根据质量守恒定律, 得到溶洞区域 I 的质量连续方程为:

$$qB = \pi r_i^2 \left(\frac{K}{\mu} \right)_{f1} \frac{\partial p_{f1}}{\partial z} \Big|_{z=h_1} - C_{cv1} \frac{dp_{cv1}}{dt} \quad (2)$$

使用相同的方式, 溶洞区域 II 界面处的压力连续方程可表示为:

$$p_{f1}(z=h_2, t) = p_{f2}(z=h_3, t) = p_{cv2} \quad (3)$$

溶洞区域 II 的质量连续方程为:

$$\pi r_i^2 \left(\frac{K}{\mu} \right)_{f1} \frac{\partial p_{f1}}{\partial z} \Big|_{z=h_2} = \pi r_i^2 \left(\frac{K}{\mu} \right)_{f2} \frac{\partial p_{f2}}{\partial z} \Big|_{z=h_3} - C_{cv2} \frac{dp_{cv2}}{dt} \quad (4)$$

由于裂缝区域 I 的高角度裂缝发育, 假设主要的流动为沿着垂直方向的一维线性流动, 由于重力方向与流动方向相反, 根据达西定律, 裂缝区域 I 任意一点的渗流速度为:

$$v_z = - \left(\frac{K}{\mu} \right)_{f1} \left(\frac{\partial p_{f1}}{\partial z} - \rho g \right) \quad (5)$$

孔洞和基质内流体以拟稳态窜流的方式流入裂缝介质, 裂缝区域 I 的质量守恒方程可写为:

$$\frac{\partial(\phi \rho)}{\partial t} + \nabla \cdot (\rho v_z) = \rho(q_{v1} + q_{m1}) \quad (6)$$

根据 Warren-Root 模型, 介质之间的窜流量可分别表示为:

$$q_{v1} = -\alpha \left(\frac{K}{\mu} \right)_{v1} (p_{v1} - p_{f1}) \quad (7)$$

$$q_{m1} = -\alpha \left(\frac{K}{\mu} \right)_{m1} (p_{m1} - p_{f1}) \quad (8)$$

结合岩石和流体的状态方程, 将(5), (7)和(8)式分别代入(6)式, 得到裂缝区域 I 中裂缝、孔洞、基质的渗流方程分别为:

$$\frac{\partial^2 p_{f1}}{\partial z^2} - 2\rho g c_f \frac{\partial p_{f1}}{\partial z} + \alpha \frac{K_{v1}}{K_{f1}} (p_{v1} - p_{f1}) + \alpha \frac{K_{m1}}{K_{f1}} (p_{m1} - p_{f1}) = \left(\frac{\phi \mu c_i}{K} \right)_{f1} \frac{\partial p_{f1}}{\partial t} \quad (9)$$

$$-\alpha \left(\frac{K}{\mu} \right)_{v1} (p_{v1} - p_{f1}) = (\phi c_i)_{v1} \frac{\partial p_{v1}}{\partial t} \quad (10)$$

$$-\alpha \left(\frac{K}{\mu} \right)_{m1} (p_{m1} - p_{f1}) = (\phi c_i)_{m1} \frac{\partial p_{m1}}{\partial t} \quad (11)$$

同理可得到裂缝区域 II 中裂缝、孔洞、基质的渗流方程分别为:

$$\frac{\partial^2 p_{f2}}{\partial z^2} - 2\rho g c_f \frac{\partial p_{f2}}{\partial z} + \alpha \frac{K_{v2}}{K_{f2}} (p_{v2} - p_{f2}) + \alpha \frac{K_{m2}}{K_{f2}} (p_{m2} - p_{f2}) = \left(\frac{\phi \mu c_i}{K} \right)_{f2} \frac{\partial p_{f2}}{\partial t} \quad (12)$$

$$-\alpha \left(\frac{K}{\mu} \right)_{v2} (p_{v2} - p_{f2}) = (\phi c_i)_{v2} \frac{\partial p_{v2}}{\partial t} \quad (13)$$

$$-\alpha \left(\frac{K}{\mu} \right)_{m2} (p_{m2} - p_{f2}) = (\phi c_i)_{m2} \frac{\partial p_{m2}}{\partial t} \quad (14)$$

模型的初始条件是:

$$p_{f1}(t=0) = p_{f2}(t=0) = p_{cv1}(t=0) = p_{cv2}(t=0) = p_i \quad (15)$$

考虑了 3 种外边界条件:

$$\begin{cases} p_{f2}(z \rightarrow \infty, t) = p_i & \text{无限大边界} \\ \frac{\partial p_{f2}}{\partial z}(z = h_4, t) = 0 & \text{封闭边界} \\ p_{f2}(z = h_4, t) = p_i & \text{定压边界} \end{cases} \quad (16)$$

1.2.2 模型求解

为了便于求解, 将(1)~(4)式和(9)~(16)式无量纲化并进行 Laplace 变换(表 1)。

表1 无量纲定义
Table1 Dimensionless definitions

变量	定义	变量	定义
无量纲压力	$p_{\text{D}} = \left(\frac{K}{\mu}\right)_{\text{fl}} \frac{\pi r_{\text{f}}(p_{\text{i}} - p_{\text{l}})}{qB}$	无量纲时间	$t_{\text{D}} = \frac{K_{\text{fl}} t}{(\phi c_{\text{l}})_{\text{fl} + \text{v1} + \text{m1}} \mu r_{\text{f}}^2}$
无量纲井筒储集系数	$C_{\text{D}} = \frac{C}{\pi(\phi c_{\text{l}})_{\text{fl} + \text{v1} + \text{m1}} r_{\text{f}}^3}$	无量纲溶洞储集系数	$C_{\text{v1D}} = \frac{C_{\text{v1}}}{\pi(\phi c_{\text{l}})_{\text{fl} + \text{v1} + \text{m1}} r_{\text{f}}^3}$
无量纲孔洞窜流系数	$\lambda_{\text{vj}} = \alpha \left(\frac{K_{\text{v}}}{K_{\text{l}}}\right) r_{\text{f}}^2$	无量纲基质窜流系数	$\lambda_{\text{mj}} = \alpha \left(\frac{K_{\text{m}}}{K_{\text{l}}}\right) r_{\text{f}}^2$
无量纲孔洞储容比	$\omega_{\text{vj}} = \frac{(\phi c_{\text{l}})_{\text{vj}}}{(\phi c_{\text{l}})_{\text{fj} + \text{vj} + \text{mj}}}$	无量纲基质储容比	$\omega_{\text{mj}} = \frac{(\phi c_{\text{l}})_{\text{mj}}}{(\phi c_{\text{l}})_{\text{fj} + \text{vj} + \text{mj}}}$
无量纲裂缝储容比	$\omega_{\text{fj}} = \frac{(\phi c_{\text{l}})_{\text{fj}}}{(\phi c_{\text{l}})_{\text{fj} + \text{vj} + \text{mj}}}$	无量纲重力系数	$G_{\text{D}} = 2\rho g c_{\text{l}} r_{\text{f}}$
无量纲长度	$z_{\text{D}} = \frac{z}{r_{\text{f}}}, z_{\text{1D}} = \frac{h_1}{r_{\text{f}}}, \dots, z_{\text{4D}} = \frac{h_4}{r_{\text{f}}}$		

得到溶洞区域的无量纲方程组为：

$$\begin{cases} \bar{p}_{\text{flD}}(z_{\text{D}} = z_{\text{1D}}) = \bar{p}_{\text{wD}}(z_{\text{D}} = 0) = \bar{p}_{\text{cv1D}} \\ -\frac{1}{s} = \frac{\partial \bar{p}_{\text{flD}}}{\partial z_{\text{D}}} \Big|_{z_{\text{D}} = z_{\text{1D}}} - C_{\text{v1D}} s \bar{p}_{\text{cv1D}} \\ \bar{p}_{\text{flD}}(z_{\text{D}} = z_{\text{2D}}) = \bar{p}_{\text{flD}}(z_{\text{D}} = z_{\text{3D}}) = \bar{p}_{\text{cv2D}} \\ \frac{\partial \bar{p}_{\text{flD}}}{\partial z_{\text{D}}} \Big|_{z_{\text{D}} = z_{\text{2D}}} = \frac{1}{M_{12}} \frac{\partial \bar{p}_{\text{flD}}}{\partial z_{\text{D}}} \Big|_{z_{\text{D}} = z_{\text{3D}}} - C_{\text{v2D}} s \bar{p}_{\text{cv2D}} \end{cases} \quad (17)$$

裂缝区域的无量纲方程组为：

$$\begin{cases} \frac{\partial^2 \bar{p}_{\text{flD}}}{\partial z_{\text{D}}^2} - G_{\text{D}} \frac{\partial \bar{p}_{\text{flD}}}{\partial z_{\text{D}}} = sD_1(s) \bar{p}_{\text{flD}} \\ \frac{\partial^2 \bar{p}_{\text{flD}}}{\partial z_{\text{D}}^2} - G_{\text{D}} \frac{\partial \bar{p}_{\text{flD}}}{\partial z_{\text{D}}} = sD_2(s) \bar{p}_{\text{flD}} \end{cases} \quad (18)$$

其中：

$$D_j(s) = \frac{M_{\text{l}j} \omega_{\text{vj}} \lambda_{\text{vj}}}{W_{\text{l}j} \lambda_{\text{vj}} + M_{\text{l}j} \omega_{\text{vj}} s} + \frac{M_{\text{l}j} \omega_{\text{mj}} \lambda_{\text{mj}}}{W_{\text{l}j} \lambda_{\text{mj}} + M_{\text{l}j} \omega_{\text{mj}} s} + \frac{M_{\text{l}j}}{W_{\text{l}j}} \omega_{\text{fj}} \quad j = 1, 2 \quad (19)$$

$$M_{\text{l}j} = \frac{\left(\frac{K}{\mu}\right)_{\text{fl}}}{\left(\frac{K}{\mu}\right)_{\text{fj}}} \quad (20)$$

$$W_{\text{l}j} = \frac{(\phi c_{\text{l}})_{\text{fl} + \text{v1} + \text{m1}}}{(\phi c_{\text{l}})_{\text{fj} + \text{vj} + \text{mj}}} \quad (21)$$

初始条件是：

$$\begin{aligned} \bar{p}_{\text{flD}}(t_{\text{D}} = 0) &= \bar{p}_{\text{flD}}(t_{\text{D}} = 0) = \\ \bar{p}_{\text{v1D}}(t_{\text{D}} = 0) &= \bar{p}_{\text{v2D}}(t_{\text{D}} = 0) = 0 \end{aligned} \quad (22)$$

外边界条件是：

$$\begin{cases} \bar{p}_{\text{flD}}(z_{\text{D}} \rightarrow \infty, s) = 0 & \text{无限大边界} \\ \frac{\partial \bar{p}_{\text{flD}}}{\partial z_{\text{D}}}(z_{\text{D}} = z_{\text{4D}}, s) = 0 & \text{封闭边界} \\ \bar{p}_{\text{flD}}(z_{\text{D}} = z_{\text{4D}}, s) = 0 & \text{定压边界} \end{cases} \quad (23)$$

求解(18)式,得到模型的通解：

$$\begin{cases} \bar{p}_{\text{flD}}(z_{\text{D}}, s) = c_1 e^{r_1 z_{\text{D}}} + c_2 e^{r_2 z_{\text{D}}} \\ \bar{p}_{\text{flD}}(z_{\text{D}}, s) = c_3 e^{r_3 z_{\text{D}}} + c_4 e^{r_4 z_{\text{D}}} \end{cases} \quad (24)$$

其中：

$$r_1 = \frac{G_{\text{D}} + \sqrt{G_{\text{D}}^2 + 4sD_1(s)}}{2} \quad (25)$$

$$r_2 = \frac{G_{\text{D}} - \sqrt{G_{\text{D}}^2 + 4sD_1(s)}}{2} \quad (26)$$

$$r_3 = \frac{G_{\text{D}} + \sqrt{G_{\text{D}}^2 + 4sD_2(s)}}{2} \quad (27)$$

$$r_4 = \frac{G_{\text{D}} - \sqrt{G_{\text{D}}^2 + 4sD_2(s)}}{2} \quad (28)$$

根据边界条件和初始条件即可确定 c_1 — c_4 ,获得模型的唯一解。

对于无限大边界：

$$\begin{cases} c_1 = \frac{e^{r_2^2 2D - r_2^2 1D}}{s \left[e^{r_1^2 2D} (r_2 - C_{v1D}s) - (r_1 - C_{v1D}s) e^{(r_1 - r_2)z_{1D} + r_2^2 2D} - \frac{(r_2 - C_{v1D}s)(r_1 - r_2) e^{r_1^2 2D}}{\frac{r_4}{M_{12}} - C_{v2D}s - r_2} \right]} \\ c_2 = -\frac{\frac{1}{s} + c_1 (r_1 - C_{v1D}s) e^{r_1^2 1D}}{(r_2 - C_{v1D}s) e^{r_2^2 1D}} \\ c_3 = 0 \\ c_4 = \frac{c_1 (r_1 - r_2) e^{r_1^2 2D}}{\left(\frac{r_4}{M_{12}} - C_{v2D}s - r_2 \right) e^{r_4^2 3D}} \end{cases} \quad (29)$$

对于封闭边界:

$$\begin{cases} c_1 = \frac{e^{r_2^2 2D}}{s(r_2 - C_{v1D}s) e^{r_2^2 1D} \left[e^{r_1^2 2D} - \frac{r_1 - C_{v1D}s}{r_2 - C_{v1D}s} e^{(r_1 - r_2)z_{1D} + r_2^2 2D} + \frac{r_1 - r_2}{X_1} \left[r_4 e^{(r_4 - r_3)z_{4D} + r_3^2 3D + r_1^2 2D} - r_3 e^{r_1^2 2D + r_4^2 3D} \right] \right]} \\ c_2 = -\frac{\frac{1}{s} + c_1 (r_1 - C_{v1D}s) e^{r_1^2 1D}}{(r_2 - C_{v1D}s) e^{r_2^2 1D}} \\ c_3 = -\frac{c_4 r_4}{r_3} e^{(r_4 - r_3)z_{4D}} \\ c_4 = \frac{c_1 r_3 (r_1 - r_2) e^{r_1^2 2D}}{X_1} \\ X_1 = r_3 \left(\frac{r_4}{M_{12}} - r_2 - C_{v2D}s \right) e^{r_4^2 3D} - r_4 \left(\frac{r_3}{M_{12}} - r_2 - C_{v2D}s \right) e^{(r_4 - r_3)z_{4D} + r_3^2 3D} \end{cases} \quad (30)$$

对于定压边界:

$$\begin{cases} c_1 = \frac{e^{r_2^2 2D}}{s(r_2 - C_{v1D}s) e^{r_2^2 1D} \left[e^{r_1^2 2D} - \frac{r_1 - C_{v1D}s}{r_2 - C_{v1D}s} e^{(r_1 - r_2)z_{1D} + r_2^2 2D} + \frac{r_1 - r_2}{X_2} \left[e^{(r_4 - r_3)z_{4D} + r_3^2 3D + r_1^2 2D} - e^{r_1^2 2D + r_4^2 3D} \right] \right]} \\ c_2 = -\frac{\frac{1}{s} + c_1 (r_1 - C_{v1D}s) e^{r_1^2 1D}}{(r_2 - C_{v1D}s) e^{r_2^2 1D}} \\ c_3 = -c_4 e^{(r_4 - r_3)z_{4D}} \\ c_4 = \frac{c_1 (r_1 - r_2) e^{r_1^2 2D}}{X_2} \\ X_2 = \left(\frac{r_4}{M_{12}} - r_2 - C_{v2D}s \right) e^{r_4^2 3D} - \left(\frac{r_3}{M_{12}} - r_2 - C_{v2D}s \right) e^{(r_4 - r_3)z_{4D} + r_3^2 3D} \end{cases} \quad (31)$$

将 c_1, c_2 代入(1)式即可得到模型的井底压力 得井底压力解^[25]:

解:

$$\bar{p}_{wD} = c_1 e^{r_1^2 1D} + c_2 e^{r_2^2 1D} \quad (32)$$

如果考虑井储和表皮效应,利用杜哈美原理获

$$\bar{p}_{wD,CS} = \frac{\bar{s}\bar{p}_{wD} + S}{s \left[1 + c_D s (\bar{s}\bar{p}_{wD} + S) \right]} \quad (33)$$

对(33)式进行 Stehfest 数值反演即可得到实空间下的井底压力解,绘制双对数图版^[26]。

2 典型图版和流动阶段

给定基础参数包括: $C_D=0.05, S=0.1, C_{v1D}=10, C_{v2D}=500, z_{1D}=1, z_{2D}=5, z_{3D}=10, z_{4D}=20, \omega_f=0.01, \omega_v=0.14, \omega_m=0.85, \lambda_v=0.001, \lambda_m=0.0001, G_D=0, M_{12}=1, W_{12}=1$, 绘制了垂向串珠结构的试井典型曲线(图2, 图3)。由图2和图3可见,试井典型曲线可划分为4个典型流动阶段:①井筒储集阶段,压力与压力导数的斜率为1。②溶洞 I 储集阶段,反映溶洞 I 内的流体流动,典型特征是压力导数曲线“下凹”且出现单位斜率直线。③溶洞 II 储集阶段,反映溶洞 II 内的流体流动,典型特征与溶洞 I 储集阶段的相同。④边界响应段,对于封闭边界,压力和压力导数曲线上翘且斜率为1;对于无限大边界,压力和压力导数曲线的斜率为1/2,反映裂缝区域线性流,这与常规无限大边界的径向流特征有差别;对于定压边界,压力导数曲线下掉。值得注意的是双溶洞模型中裂缝区域的流动特征被溶洞区域掩盖了。为了阐述完整的流动阶段,给出单溶洞单裂缝区域串联的试井典型曲线。裂缝区域主要包含3个流动阶

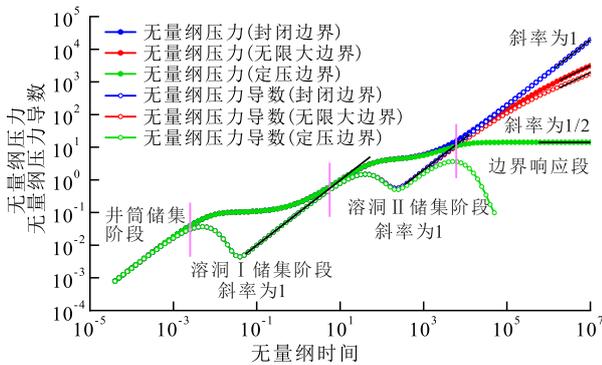


图2 双溶洞双裂缝区域模型试井典型图版
Fig.2 Type curves of two-cave-two-region model

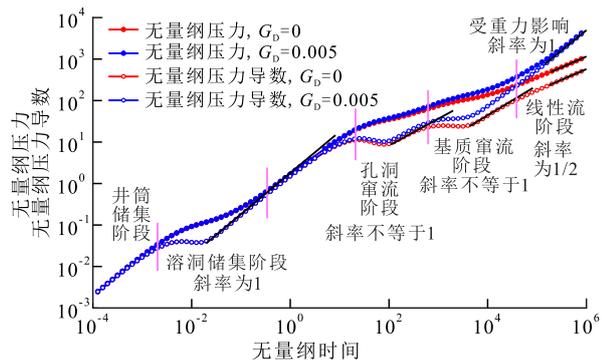


图3 单溶洞单裂缝区域模型试井典型图版
Fig.3 Type curves of single-cave-single-region model

段(图3):①孔洞窜流阶段,压力导数曲线“下凹”。②基质窜流阶段,压力导数曲线“下凹”。③线性流阶段,压力和压力导数曲线的斜率为1/2。与溶洞介质引起的压力导数曲线“下凹”不同,常规多重介质之间窜流引起的压力导数曲线“下凹”的斜率往往不等于1。这为识别溶洞和常规的多重介质提供了理论依据。

另外,分析了重力对试井典型曲线的影响。与封闭边界相似,重力影响下试井典型曲线在后期出现斜率为1的压力和压力导数曲线(图3)。因此,开发串珠状缝洞油藏应考虑重力对流动的影响。顺北油田采取的“注水替油”措施减轻了重力的影响,已经取得了较好的开发效果^[27]。

3 参数敏感性

给定基础参数 $C_D=0.05, S=0.1, C_{v1D}=10, z_{1D}=1, z_{2D}=5, z_{3D}=10, z_{4D}=20, \omega_f=0.01, \omega_v=0.14, \omega_m=0.85, \lambda_v=0.001, \lambda_m=0.0001, G_D=0, M_{12}=1, W_{12}=1$, 分析参数敏感性。

3.1 无量纲溶洞储集系数

设置3组不同无量纲溶洞储集系数,即 C_{v2D} 分别为300,500,700。由图4可见,无量纲溶洞储集系数主要影响溶洞储集阶段的宽度和深度,无量纲溶洞储集系数越大,下凹的宽度和深度越大。

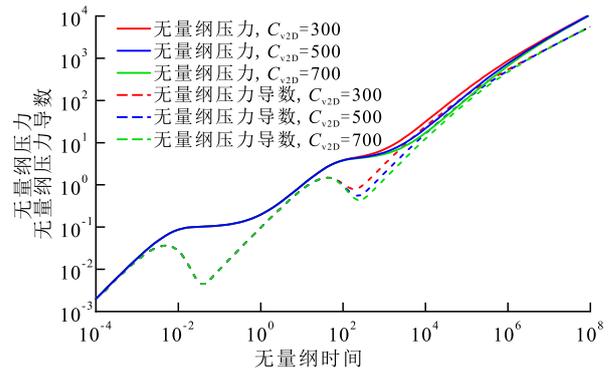


图4 无量纲溶洞储集系数对试井典型曲线的影响
Fig.4 Effect of dimensionless cave storage coefficient on type curves

3.2 裂缝长度

保持溶洞 I 的高度不变,改变 Z_{2D} 来表示不同的裂缝长度,即 Z_{2D} 分别为3,5,7。由图5可见,裂缝长度主要影响溶洞 II 储集阶段的出现时间,裂缝长度越大,溶洞 II 储集阶段出现时间越晚,因为压力波需要更长的时间穿过裂缝区域。

3.3 无量纲重力系数

设置3组不同无量纲重力系数,即 G_D 分别为0,

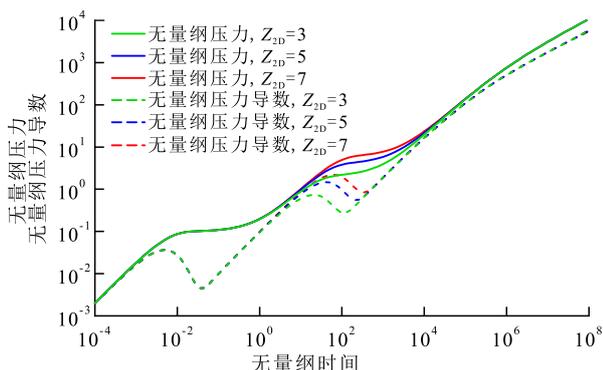


图5 裂缝长度对试井典型曲线的影响
Fig.5 Effect of fracture length on type curves

0.01, 0.001。由图6可见,重力影响下试井典型曲线在后期会出现斜率为1的压力和压力导数曲线,曲线上翘的时间与无量纲重力系数有关,无量纲重力系数越大,曲线上翘的时间越早。因此,重力可以视为封闭边界,重力的大小相当于封闭边界的距离。

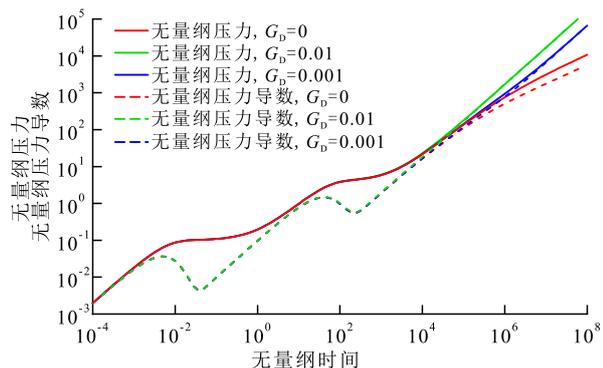


图6 无量纲重力系数对试井典型曲线的影响
Fig.6 Effect of dimensionless gravity coefficient on type curves

3.4 窜流系数

设置3组不同的基质窜流系数,即 λ_m 分别为0.000 1, 0.000 2, 0.000 3,来分析其对试井典型曲线的影响。由图7可见,与常规认识一致,窜流系数主要影响窜流阶段出现的时间,窜流系数越大,窜流阶段出现的时间越早。

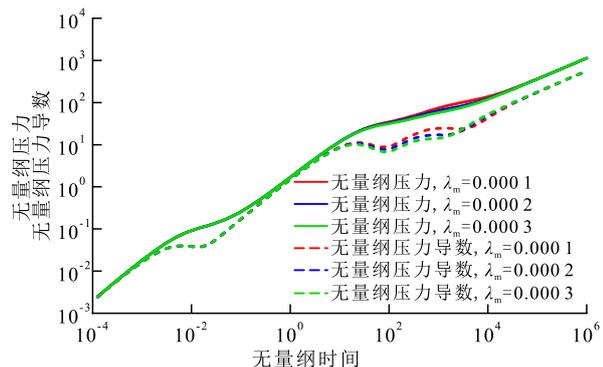


图7 窜流系数对试井典型曲线的影响
Fig.7 Effect of interporosity flow coefficient on type curves

3.5 储容比

改变孔洞储容比和基质储容比,即 ω_v 分别为0.10, 0.14, 0.18, ω_m 分别为0.81, 0.85, 0.89,来分析其对试井典型曲线的影响。由图8可见,储容比主要影响窜流阶段下凹的深度,储容比越大,窜流阶段下凹的深度越大。

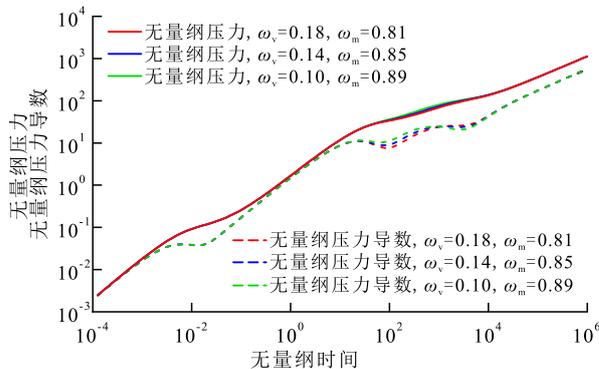


图8 储容比对试井典型曲线的影响
Fig.8 Effect of storativity ratio on type curves

4 实例应用

顺北油田是塔里木盆地新发现的断控岩溶油藏,构造破碎带控制着油藏的主要储量,地震剖面显示构造破碎带发育大量的串珠状缝洞体,钻井过程中经常发生钻杆放空和泥浆漏失的情况。现有试井分析方法解释的参数很难与地质情况吻合,为了阐述新建模型的应用,选取顺北油田典型井进行解释。

顺北Y井钻入垂向串珠状缝洞体(图9),日产

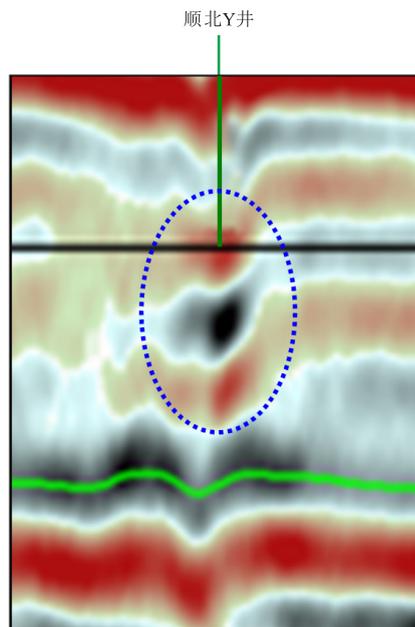


图9 顺北Y井地震剖面
Fig.9 Seismic profile of Well Shunbei Y

量为 65 m³/d, 压力恢复测试时长为 255 h, 实测试井曲线出现了溶洞储集阶段和窜流阶段, 该井基础参数包括: 井径为 0.1 m, 储层有效孔隙度为 10%, 综合压缩系数为 0.000 836 MPa⁻¹, 地层原油体积系数为 1.114, 地层原油黏度为 2.242 mPa·s。使用提出的模型对该井进行解释, 理论双对数曲线与实测数据的拟合效果较好(图 10), 表 2 列出了试井解释结果, 油藏的平均渗透率为 796 mD, 渗透性较好。此外, 还获得了常规试井方法无法解释的溶洞等效半径为 52 m, 溶洞 I 等效高度为 2.6 m。

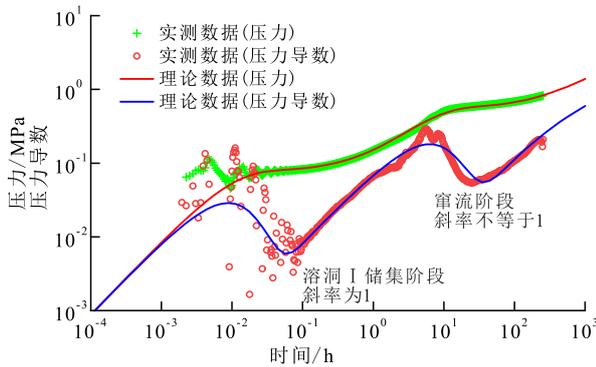


图 10 顺北 Y 井压力和压力导数拟合结果

Fig.10 Fitting results of pressure and its derivative of Well Shunbei Y

表 2 顺北 Y 井解释结果

Table2 Interpretation results of Well Shunbei Y

油藏参数	数值	油藏参数	数值
$C/(m^3 \cdot MPa^{-1})$	0.35	S	5.26
$C_{v1}/(m^3 \cdot MPa^{-1})$	35.1	h_1/m	2.6
K/mD	796	r_f/m	52
ω_f	0.006	λ_v	0.000 92

5 结论

详细分析了 3 种外边界条件下垂向串珠状缝洞体的试井典型曲线, 其典型流动阶段包括溶洞储集阶段和裂缝区域线性流阶段。

溶洞储集阶段和窜流阶段的压力导数曲线都会出现“下凹”, 溶洞储集阶段的典型特征是压力导数曲线的斜率为 1, 而窜流阶段压力导数曲线的斜率不等于 1, 这为通过试井技术识别溶洞和常规的多重介质提供了理论依据。

在串珠状缝洞体中, 重力影响下试井典型曲线会出现类似封闭边界的压力响应特征, 开发过程中应采取措施减轻重力的影响以提高串珠状缝洞型油藏的开发效果, 顺北油田实施的“注水替油”措施已取得不错的效果。

顺北油田发育垂向串珠状缝洞体, 应用新建模型解释了一口典型井, 除能解释常规的储层参数外, 也能反演现有试井方法无法解释的溶洞等效半径、溶洞等效高度等参数。

符号解释

- B ——原油体积系数;
- c_1, c_2, c_3, c_4 ——常数;
- c_f ——流体压缩系数, Pa⁻¹;
- c_i ——综合压缩系数, Pa⁻¹;
- 下标 cv——溶洞;
- C ——井筒储集系数, m³/Pa;
- C_{cv1} ——溶洞区域 I 储集系数, m³/MPa;
- C_{cv2} ——溶洞区域 II 储集系数, m³/MPa;
- C_{vi} ——溶洞区域 i 储集系数, m³/Pa;
- 下标 D——无量纲;
- $D_j(s)$ ——裂缝区域 j 窜流函数;
- 下标 f——裂缝介质;
- g ——重力加速度, m/s², 其值为 9.8;
- G_D ——无量纲重力系数;
- h ——溶洞等效高度, m;
- h_1 ——溶洞区域 I 底深, 即溶洞 I 等效高度, m;
- h_2 ——裂缝区域 I 底深, m;
- h_3 ——溶洞区域 II 底深, m;
- h_4 ——裂缝区域 II 底深, m;
- 下标 i ——溶洞序号, 分别为 1 和 2;
- 下标 j ——裂缝序号, 分别为 1 和 2;
- K ——渗透率, D;
- K_f, K_m, K_v ——裂缝区域 j 的裂缝、基质和孔洞渗透率, D;
- D ;
- K_f, K_m, K_v ——裂缝、基质和孔洞渗透率, D;
- 下标 m——基质介质;
- M_j ——裂缝区域 I 和裂缝区域 j 的流度比;
- p_{cv1} ——第 i 个溶洞区域压力, Pa;
- p_{cv1} ——溶洞区域 I 压力, Pa;
- p_{cv2} ——溶洞区域 II 压力, Pa;
- p_f ——裂缝压力, Pa;
- p_{fj} ——裂缝区域 j 裂缝介质压力, Pa;
- p_{f1} ——裂缝区域 I 裂缝介质压力, Pa;
- p_{f2} ——裂缝区域 II 裂缝介质压力, Pa;
- p_i ——原始地层压力, Pa;
- p_{m1} ——裂缝区域 I 基质介质压力, Pa;
- p_{m2} ——裂缝区域 II 基质介质压力, Pa;
- p_{v1} ——裂缝区域 I 孔洞介质压力, Pa;
- p_{v2} ——裂缝区域 II 孔洞介质压力, Pa;
- p_{wf} ——井底流压, Pa;
- $p_{wf,CS}$ ——考虑井储表皮的井底流压, Pa;
- q ——油井产量, m³/s;

q_{mi} ——裂缝区域 I 基质向裂缝的窜流量, m^3/s ;
 q_{vi} ——裂缝区域 I 孔洞向裂缝的窜流量, m^3/s ;
 r ——径向坐标系;
 r_1, r_2, r_3, r_4 ——常数;
 r_f ——溶洞等效半径, m ;
 r_w ——井筒半径, m ;
 s ——Laplace 变量, 无量纲;
 S ——表皮系数;
 t ——生产时间, s ;
 下标 v ——孔洞介质;
 v_z ——裂缝区域任意一点的渗流速度, m/s ;
 W_{ij} ——裂缝区域 I 和裂缝区域 j 的储容比;
 X_1, X_2 ——常数;
 z ——纵向距离, m ;
 $z_D, z_{1D}, z_{2D}, z_{3D}, z_{4D}$ ——无量纲纵向距离;
 α ——形状因子;
 λ_m, λ_v ——基质、孔洞窜流系数;
 $\lambda_{mj}, \lambda_{vj}$ ——裂缝区域 j 的基质、孔洞窜流系数;
 μ ——流体黏度, $Pa \cdot s$;
 $\mu_{ij}, \mu_{mj}, \mu_{vj}$ ——裂缝区域 j 的裂缝、基质和孔洞流体黏度, $Pa \cdot s$;
 ρ ——流体密度, kg/m^3 ;
 ϕ ——孔隙度;
 $\phi_{ij}, \phi_{mj}, \phi_{vj}$ ——裂缝区域 j 的裂缝、基质和孔洞孔隙度;
 $\omega_f, \omega_m, \omega_v$ ——裂缝、基质和孔洞储容比;
 $\omega_{ij}, \omega_{mj}, \omega_{vj}$ ——裂缝区域 j 的裂缝、基质、孔洞储容比。

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编辑 刘北羿